4.2 Energy and Power

Definition 4.12. For a signal g(t), the instantaneous power p(t) dissipated in the 1- Ω resister is $p_g(t) = |g(t)|^2$ regardless of whether g(t) represents a voltage or a current. To emphasize the fact that this power is based upon unity resistance, it is often referred to as the **normalized** (instantaneous) power.

Definition 4.13. The total (normalized) **energy** of a signal g(t) is given by

$$E_g = \int_{-\infty}^{+\infty} p_g(t)dt = \int_{-\infty}^{+\infty} |g(t)|^2 dt = \lim_{T \to \infty} \int_{-T}^{T} |g(t)|^2 dt.$$

4.14. By the **Parseval's theorem** discussed in 2.43, we have

$$E_g = \int_{-\infty}^{\infty} |g(t)|^2 dt = \int_{-\infty}^{\infty} |G(f)|^2 df.$$

Definition 4.15. The average (normalized) **power** of a signal g(t) is given by

$$P_g = \lim_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} |g(t)|^2 dt = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} |g(t)|^2 dt.$$

Definition 4.16. To simplify the notation, there are two operators that used angle brackets to define two frequently-used integrals:

(a) The "time-average" operator:

$$\langle g \rangle \equiv \langle g(t) \rangle \equiv \lim_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} g(t) dt = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} g(t) dt$$
 (42)

(b) The **inner-product** operator:

$$\langle g_1, g_2 \rangle \equiv \langle g_1(t), g_2(t) \rangle = \int_{-\infty}^{\infty} g_1(t) g_2^*(t) dt$$
 (43)

4.17. Using the above definition, we may write

•
$$E_g = \langle g, g \rangle = \langle G, G \rangle$$
 where $G = \mathcal{F} \{g\}$

$$\bullet P_g = \langle |g|^2 \rangle$$

= a < 9(t)> + b < 9,1t>>

- Parseval's theorem: $\langle g_1, g_2 \rangle = \langle G_1, G_2 \rangle$ where $G_1 = \mathcal{F} \{g_1\}$ and $G_2 = \mathcal{F} \{g_2\}$
- **4.18.** Time-Averaging over Periodic Signal: For **periodic** signal g(t) with period T_0 , the time-average operation in (42) can be simplified to

$$\langle g \rangle = \frac{1}{T_0} \int_{T_0} g(t) dt$$

where the integration is performed over a period of q.

where the integration is performed over a period of g.

Example 4.19. $\langle \cos(2\pi f_0 t + \theta) \rangle = \begin{cases} 0, & \forall \theta \neq 0, \\ \cos(\theta), & \forall \theta \neq 0, \end{cases}$ Similarly, $\langle \sin(2\pi f_0 t + \theta) \rangle = \begin{cases} 0, & \forall \theta \neq 0, \\ \sin(\theta), & \forall \theta \neq 0, \end{cases}$ Example 4.20. $\langle \cos^2(2\pi f_0 t + \theta) \rangle = \langle \frac{1}{2} + \frac{1}{2}\cos(2\pi f_0 t + 2\theta) \rangle$ $\cos^2 x = \frac{1}{2}(1+\cos(2\pi t))$ $= \begin{cases} 1/2, & \forall \theta \neq 0, \\ \cos^2 \theta, & \forall \theta \neq 0, \\ \cos^2 \theta, & \forall \theta \neq 0, \end{cases}$

Example 4.21. $\langle e^{j(2\pi f_0 t + \theta)} \rangle = \langle \cos(2\pi f_0 t + \theta) + j \sin(2\pi f_0 t + \theta) \rangle$ $= \begin{cases} 0, & \text{if } 0 < 0, \\ \text{if } 0 < 0, \\ \text{if } 0 < 0, \end{cases}$

Example 4.22. Suppose $g(t) = ce^{j2\pi f_0 t}$ for some (possibly complex-valued) constant c and (real-valued) frequency f_0 . Find P_q .

$$P_g = \langle |g|^2 \rangle = \langle |ce^{j2\pi/6t}|^2 \rangle = \langle |c|^2 |e^{j2\pi/6t}|^2 \rangle = |c|^2$$

4.23. When the signal g(t) can be expressed in the form $g(t) = \sum_{k} c_k e^{j2\pi f_k t}$ and the f_k are distinct, then its (average) power can be calculated from

$$P_g = \sum_k \left| c_k \right|^2$$

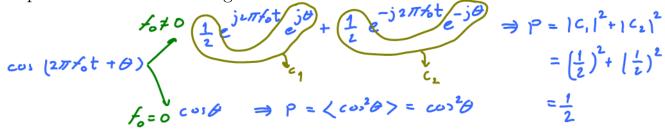
Example 4.24. Suppose
$$g(t)=2e^{j6\pi t}+3e^{j8\pi t}$$
. Find P_g .

Example 4.25. Suppose $g(t) = 2e^{j6\pi t} + 3e^{j6\pi t}$. Find P_g .

$$j_{2}\pi(3)$$
t
g(t) = 5e
 $P_{g} = 151^{2} = 25$

Example 4.26. Suppose $g(t) = \cos(2\pi f_0 t + \theta)$. Find P_g .

Here, there are several ways to calculate P_g . We can simply use Example 4.20. Alternatively, we can first decompose the cosine into complex exponential functions using the Euler's formula:



4.27. The (average) power of a sinusoidal signal $g(t) = A\cos(2\pi f_0 t + \theta)$ is

$$P_g = \begin{cases} \frac{1}{2}|A|^2, & f_0 \neq 0, \\ |A|^2 \cos^2 \theta, & f_0 = 0. \end{cases}$$

This property means any sinusoid with nonzero frequency can be written in the form

$$g(t) = \sqrt{2P_g}\cos(2\pi f_0 t + \theta).$$

4.28. Extension of 4.27: Consider sinusoids $A_k \cos(2\pi f_k t + \theta_k)$ whose frequencies are positive and distinct. The (average) power of their sum

$$g(t) = \sum_{k} A_k \cos(2\pi f_k t + \theta_k)$$

is

$$P_g = \frac{1}{2} \sum_{k} |A_k|^2.$$



Example 4.29. Suppose
$$g(t) = 2\cos(2\pi\sqrt{3}t) + 4\cos(2\pi\sqrt{5}t)$$
. Find P_g .
$$P_g = |A_1|^2 + |A_2|^2 = \frac{1}{2}(2^2 + 7^2) = \frac{1}{2}(4 + 16) = 10$$

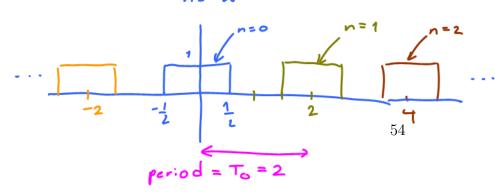
Example 4.30. Suppose $g(t) = 3\cos(2t) + 4\cos(2t - 30^{\circ}) + 5\sin(3t)$. Find P_q .

$$P_9 \approx \frac{1}{2} \left(677^2 + 5^2 \right) \approx 35.4$$

4.31. For **periodic** signal q(t) with period T_0 , there is also no need to carry out the limiting operation to find its (average) power P_g . We only need to find an average carried out over a single period:

$$P_g = \frac{1}{T_0} \int_{T_0} |g(t)|^2 dt.$$

Example 4.32.



$$rect(t) \qquad \begin{array}{c|c} 1 \\ \hline \\ -\frac{1}{2} & \frac{1}{2} \end{array}$$

$$P_{g} = \frac{1}{T_{o}} \int |g|^{2} t_{o}| dt$$

$$= \frac{1}{2} \int_{\frac{1}{2}0}^{\infty} |g|^{2} t_{o}| dt$$

$$= \frac{1}{2} \int_{\frac{1}{2}0}^{\infty} |g|^{2} dt$$

- **4.33.** When the Fourier series expansion (to be reviewed in Section 4.3) of the signal is available, it is easy to calculate its power:
- (a) When the corresponding Fourier series expansion $g(t) = \sum_{k=-\infty}^{\infty} c_k e^{j2\pi k f_0 t}$ is known,

$$P_g = \sum_{k=-\infty}^{\infty} |c_k|^2.$$

(b) When the signal g(t) is real-valued and its (compact) trigonometric Fourier series expansion $g(t) = c_0 + 2\sum_{k=1}^{\infty} |c_k| \cos(2\pi k f_0 t + \angle \phi_k)$ is known,

$$P_g = c_0^2 + 2\sum_{k=1}^{\infty} |c_k|^2$$
.

Definition 4.34. Based on Definitions 4.13 and 4.15, we can define three distinct classes of signals:

- (a) If E_q is finite and nonzero, g is referred to as an energy signal.
- (b) If P_q is finite and nonzero, g is referred to as a **power signal**.
- (c) Some signals¹⁷ are neither energy nor power signals.
 - Note that the power signal has infinite energy and an energy signal has zero average power; thus the two categories are disjoint.

 $g(t) = N \times \text{rect}(t)$ $= \int_{-\frac{1}{2}}^{\infty} |g(t)|^2 dt = h^2$ $= \int_{-\frac{1}{2}}^{\infty} |g(t)|^2 dt$ **Example 4.35.** Rectangular pulse

¹⁷Consider $g(t) = t^{-1/4} 1_{[t_0,\infty)}(t)$, with $t_0 > 0$.

Example 4.36. Sinc pulse

Example 4.37. For $\alpha > 0$, $g(t) = Ae^{-\alpha t}1_{[0,\infty)}(t)$ is an energy signal with $E_g = |A|^2/2\alpha$.

Example 4.38. The rotating phasor signal $g(t) = ce^{j(2\pi f_0 t + \theta)}$ is a power signal with $P_g = |c|^2$.

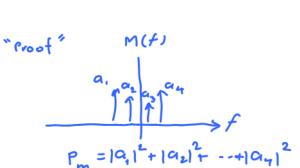
Example 4.39. The sinusoidal signal $g(t) = A\cos(2\pi f_0 t + \theta)$ is a power signal with $P_g = |A|^2/2$.

4.40. Consider the transmitted signal

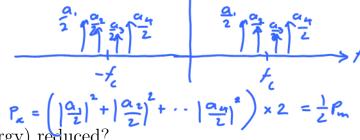
$$x(t) = m(t)\cos(2\pi f_c t + \theta)$$

in DSB-SC modulation. Suppose $M(f - f_c)$ and $M(f + f_c)$ do not overlap (in the frequency domain).

(a) If m(t) is a power signal with power P_m , then the average transmitted power is



• Q: Why is the power (or energy) reduced?



- Remark: When $x(t) = \sqrt{2}m(t)\cos(2\pi f_c t + \theta)$ (with no overlapping between $M(f f_c)$ and $M(f + f_c)$), we have $P_x = P_m$.
- (b) If m(t) is an energy signal with energy E_m , then the transmitted energy is

$$E_x = \frac{1}{2}E_m.$$

Example 4.41. Suppose $m(t) = \cos(2\pi f_c t)$. Find the average power in $x(t) = m(t)\cos(2\pi f_c t)$.